

## Improving retrieval of volcanic sulfur dioxide from backscattered UV satellite observations

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[1] Existing algorithms that use satellite measurements of solar backscattered ultraviolet (BUV) radiances to retrieve sulfur dioxide (SO<sub>2</sub>) vertical columns underestimate the large SO<sub>2</sub> amounts encountered in fresh volcanic eruption clouds. To eliminate this underestimation we have developed a new technique, named the Iterative Spectral Fitting (ISF) algorithm, for accurate retrieval of SO<sub>2</sub> vertical columns in the full range of volcanic emissions. The ISF algorithm is applied to Ozone Monitoring Instrument (OMI) BUV measurements of the Sierra Negra eruption (Galápagos Islands, Ecuador) in October 2005. The results represent major improvements over the operational OMI SO<sub>2</sub> products. Based on the ISF data, we report the largest SO<sub>2</sub> vertical column amount (>1000 Dobson Units (DU), where 1 DU =  $2.69 \times 10^{16}$  molecules/cm<sup>2</sup>) ever observed by a space borne instrument, implying that very high concentrations of SO<sub>2</sub> can occur in the lower troposphere during effusive eruptions. **Citation:** Yang, K., N. A. Krotkov, A. J. Krueger, S. A. Carn, P. K. Bhartia, and P. F. Levelt (2009), Improving retrieval of volcanic sulfur dioxide from backscattered UV satellite observations, *Geophys. Res. Lett.*, 36, L03102, doi:10.1029/2008GL036036.

### 1. Introduction

[2] Volcanic activity releases various gases and sometimes ash into the atmosphere. These emissions vary in composition, amount, duration, and height, depending on the volcano and the type and intensity of activity. Sulfur dioxide (SO<sub>2</sub>) is one of the dominant gases released by volcanic eruptions. Its subsequent dispersion, transport, and conversion to sulfate aerosol may lead to acid rain and air pollution. Large explosive eruptions can inject tremendous amounts of SO<sub>2</sub> and ash into the upper troposphere and stratosphere in a short period of time. While volcanic ash presents severe hazards to aviation safety, it usually does not lead to a significant reduction of solar radiation received at the Earth surface due to rapid fallout. On the other hand, massive amounts of SO<sub>2</sub> that reach the stratosphere are eventually converted to sulfate aerosol which can have

lasting impacts, such as reduced global surface temperatures and depletion of stratospheric ozone [McCormick *et al.*, 1995]. Even non-explosive eruptions that emit little ash might have severe regional climate effects if the amount of SO<sub>2</sub> released into the troposphere is large enough, as exemplified by the 1783 Laki (Iceland) eruption that lasted 8–9 months and led to large scale crop failure and livestock loss and the deaths of 24% of the population in Iceland [Sigurdsson, 1982].

[3] Ever since the first sighting of SO<sub>2</sub> clouds from the 1982 El Chichón eruption by the Nimbus 7 Total Ozone Mapping Spectrometer (TOMS) [Krueger, 1983], satellite measurements have played an indispensable role in assessing and quantifying volcanic emissions, which would be otherwise impossible to monitor globally due to the large size of volcanic eruption clouds, the sporadic nature of volcanic activity and the remote locations of many volcanoes. A unique long-term record of volcanic SO<sub>2</sub> emissions has been compiled from the series of TOMS observations [Carn *et al.*, 2003]. This record is being continued and improved using observations from the Ozone Monitoring Instrument (OMI) [Levelt *et al.*, 2006a] on NASA's Aura platform and supplemented by other backscattered ultraviolet (BUV) measurements from the Global Ozone Monitoring Experiment (GOME) [Burrows *et al.*, 1999], Scanning Imaging Absorption Spectrometer for Atmospheric Cartography (SCIAMACHY) [Bovensmann *et al.*, 1999], and GOME-2 [Munro *et al.*, 2006]. A number of algorithms [see Yang *et al.*, 2007, and references therein] have been developed to retrieve SO<sub>2</sub> from discrete or hyper-spectral BUV measurements in the spectral range 308 – 360 nm. The SO<sub>2</sub> vertical columns derived using these algorithms are impacted by the nonlinear effect [Yang *et al.*, 2007], depending on the SO<sub>2</sub> amount and the wavelengths used.

[4] In this paper, we describe OMI operational data (section 2) and a new iterative spectral fitting (ISF) algorithm that performs simultaneous retrievals of ozone and SO<sub>2</sub> with a proper treatment of the nonlinear effect (section 3). We compare operational and ISF results for an effusive eruption of Sierra Negra volcano (Galápagos Islands) in October 2005 to demonstrate that substantial improvements are achieved using the ISF algorithm (section 4).

### 2. OMI Data

[5] OMI is a Dutch-Finnish built UV/visible spectrometer that has been operating successfully on NASA's Aura spacecraft since its launch in July 2004 [Levelt *et al.*, 2006a]. It has maintained a stable radiometric calibration while making measurements of direct and backscattered solar radiation with hyper-spectral resolution, covering the

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globe daily (when in global mode) at a high nadir spatial resolution. The hyper-spectral measurements have facilitated an accurate wavelength calibration [Caspar and Chance, 1997; Liu *et al.*, 2005] and the characterization of the spectral response of the instrument [Liu *et al.*, 2005]. The unique combination of OMI observation characteristics has been exploited [Levelt *et al.*, 2006b, and references therein] to improve determination of effective cloud pressures, measurements of aerosol properties, and retrievals of a number of trace gases relevant to climate and air quality, including ozone, NO<sub>2</sub>, HCHO, BrO, and SO<sub>2</sub>.

[6] For total ozone and SO<sub>2</sub> retrievals, we have used data from the OMI UV-2 channel, which has a nadir footprint of  $13 \times 24 \text{ km}^2$  in the spectral range 310 – 380 nm at a sampling rate of  $\sim 0.15 \text{ nm/pixel}$  and a spectral resolution of  $\sim 0.42 \text{ nm}$ . This spectral range covers the ozone Huggins band and a large portion of the strong and weak absorption features of the SO<sub>2</sub> band. The measured spectra contain well-distinguished spectral structures of the Ring effect resulting from Rotational Raman Scattering (RRS) [see Joiner and Vasilkov, 2006, and references therein] and absorptions of constituent trace gases because of the high sampling rate and narrow spectral resolution of OMI's UV-2 channel, enabling optimal selection of measurements at specific wavelengths to maximize detection sensitivity to low SO<sub>2</sub> loading. Consequently OMI observations have provided unprecedented measurements of SO<sub>2</sub> emissions from anthropogenic and natural sources, and a 150-fold improvement in sensitivity over TOMS SO<sub>2</sub> measurements [Krotkov *et al.*, 2006].

### 3. Iterative Spectral Fit Algorithm

[7] Two retrieval techniques, the band residual difference (BRD) method [Krotkov *et al.*, 2006] for boundary layer SO<sub>2</sub> and the Linear Fit (LF) algorithm [Yang *et al.*, 2007] for volcanic SO<sub>2</sub>, are being used operationally to produce SO<sub>2</sub> vertical columns from OMI UV-2 measurements. The LF algorithm is designed to work with both discrete-wavelength and hyper-spectral measurements and the BRD method can also be extended to include measurements at more wavelengths; thus far they have been applied to measurements at a small set of wavelengths only. Although both techniques have not fully exploited the hyper-spectral capability of OMI, they have successfully measured a wide range of SO<sub>2</sub> loadings encountered in global observations [Carn *et al.*, 2007; Yang *et al.*, 2007; Krotkov *et al.*, 2008], and are fast enough to generate SO<sub>2</sub> products in near real time for aviation hazard mitigation [Carn *et al.*, 2008]. Limitations in retrieval accuracy for large eruptions are expected because these techniques assume an approximate linear relationship between SO<sub>2</sub> column and the residuals at a set of wavelengths [Yang *et al.*, 2007]. Residuals here are calculated as N-value ( $N = -100 \log_{10} I$ , where  $I$  is the sun normalized radiance) differences between the measurements and forward model simulations with an initial guess of zero SO<sub>2</sub>. The linear approximation is justified when the amount of SO<sub>2</sub> is small, because the mean path traversing the atmosphere by the measured BUV photons is almost identical to that of the BUV photons in a SO<sub>2</sub>-free atmosphere. However with a larger SO<sub>2</sub> amount, the stronger absorption suppresses more subsequent photon scatterings, reducing

the relative contribution to the measured radiance by photons that have passed the SO<sub>2</sub> medium, leading to an underestimation of the actual SO<sub>2</sub> value by the BRD and LF methods. The LF method has a wider valid range than the BRD algorithm primarily because the former uses longer wavelength measurements, where SO<sub>2</sub> absorption is weaker and the linear approximation holds for SO<sub>2</sub> loadings up to  $\sim 100 \text{ DU}$  [Yang *et al.*, 2007]. For larger SO<sub>2</sub> loadings, the non-linear response of the N-value residuals to the change in SO<sub>2</sub> becomes the main limiting factor of the operational OMI SO<sub>2</sub> algorithms.

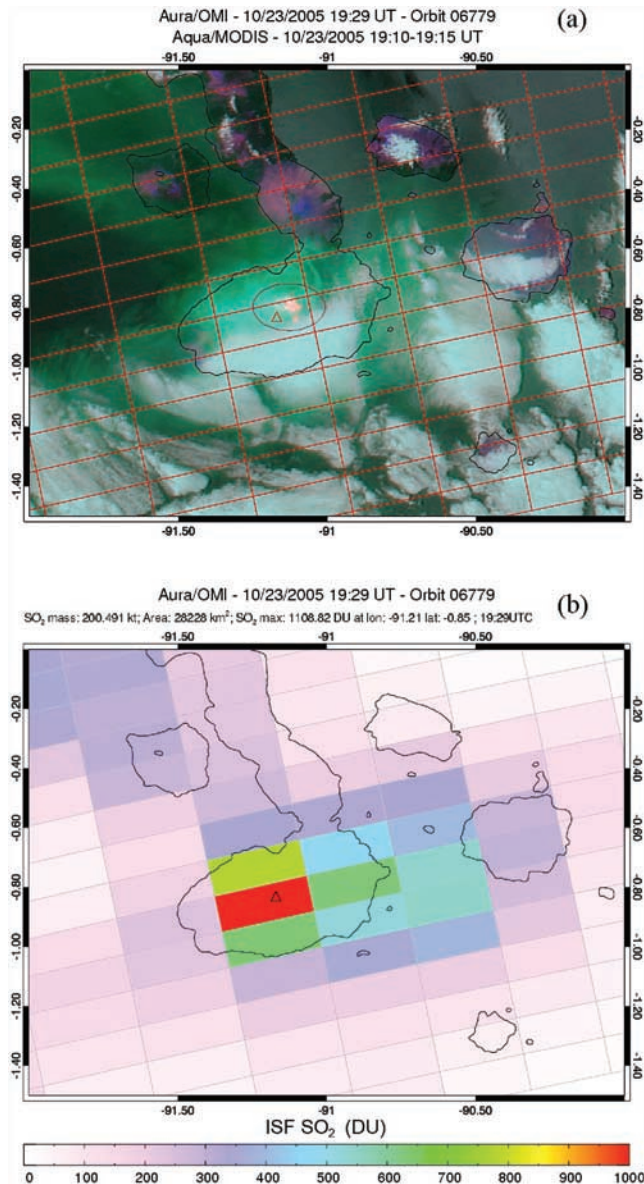
[8] To overcome this limitation due to the non-linear effect, the LF equation can be solved iteratively with an updated linearization point and weighting functions at each step. Algebraically, the LF equation for a wavelength band is expressed as

$$N_{obs} - N_i = \Delta\Omega \left. \frac{\partial N}{\partial \Omega} \right|_{\Omega=\Omega_i} + \Delta\Xi \left. \frac{\partial N}{\partial \Xi} \right|_{\Xi=\Xi_i} + p_{rrs} \sigma_i^{rrs} + \left( \Delta F_c + \sum_{k=1}^{n=2} c_k (\lambda - \lambda_0)^k \right) \left. \frac{\partial N}{\partial F_c} \right|_{F_c=F_{ci}} + \varepsilon, \quad (1)$$

where  $N_{obs}$  is the OMI measured N-value,  $N_i$  is the computed N-value at the  $i$ th iteration based on a forward model, and  $\varepsilon$  is the total error for the wavelength band. Starting with an initial guess of ozone ( $\Omega_0$ ) based on the seasonal zonal mean value at the location of the observation and zero SO<sub>2</sub> ( $\Xi_0 = 0$ ), the cloud fraction parameter ( $F_c$ ) of a Mixed Lambert-Equivalent Reflectivity (MLER) model [Ahmad *et al.*, 2004] is determined at  $\lambda_0 = 333 \text{ nm}$  (where ozone and SO<sub>2</sub> absorption cross sections and the RRS effect are small), with MLER cloud pressure determined by the effective cloud pressure from the OMI RRS (OMCLDRR) product [Joiner and Vasilkov, 2006]. Note that when the effective scene reflectivity is either low ( $< 0.15$ , usually under partial cloud conditions or in the presence of non-absorbing aerosols) or very high ( $> 0.8$ , usually when the surface is covered by snow/ice), the Lambert Equivalent Reflectivity (LER) model is used (i.e.,  $F_c$  is replaced by LER) so that the forward model can better approximate the mean path of the BUV photons measured by the satellite instrument. The linear set of equations (1) for each wavelength can be solved by least squares fitting of residuals ( $N_{obs} - N_i$ ) to the weighting functions ( $\frac{\partial N}{\partial \Omega}$ ,  $\frac{\partial N}{\partial \Xi}$ ,  $\frac{\partial N}{\partial F_c}$ ,  $(\lambda - \lambda_0) \frac{\partial N}{\partial F_c}$ ,  $(\lambda - \lambda_0)^2 \frac{\partial N}{\partial F_c}$ ) and the RRS spectrum ( $\sigma_i^{rrs}$ ) [Sioris and Evans, 1999] to obtain state vector adjustments:  $\Delta F_c$ , polynomial coefficients ( $c_1$  and  $c_2$ ), RRS scaling parameter ( $p_{rrs}$ ), and the linearization point for the next iteration ( $\Omega_{i+1} = \Omega_i + \Delta\Omega$ ,  $\Xi_{i+1} = \Xi_i + \Delta\Xi$ ). The weighting functions and the RRS spectrum are updated at each iterative step. The final solution of ozone and SO<sub>2</sub> columns is achieved when the iteration converges, i.e., the changes in ozone and SO<sub>2</sub> between successive iterations are below certain thresholds. The error due to the non-linear effect, which manifests itself in the higher order terms of  $\Delta\Omega$  and  $\Delta\Xi$  (absorbed into  $\varepsilon$  in equation (1)), becomes negligible as the iterative process converges.

[9] The TOMRAD forward model, which implements Dave's iterative solution to the vector radiative transfer equation [Dave, 1964] with a pseudo-spherical correction [Caudill *et al.*, 1997], is used to calculate the upwelling





**Figure 1.** The MODIS instrument on the Aqua satellite observed the Sierra Negra eruption about 15 minutes before OMI on October 23, 2005. (a) A color composite of MODIS bands 7 (2114.1 nm, Red), 2 (856.5 nm, Green), and 1 (645.5 nm, Blue) over the Galápagos Islands and (b) the OMI ISF SO<sub>2</sub> map of the same area. OMI pixel footprints are overlaid on both Figures 1a and 1b. Sierra Negra volcano (denoted by the symbol  $\Delta$  in Figures 1a and 1b) is located at 91.17°W and 0.83°S and rises to an elevation of 1.1 km. Oval outline in Figure 1a encloses the lava flows (orange hot spot) seen by MODIS.

radiance ( $I$ ) at each wavelength and the corresponding weighting functions. Aerosols are not included in the forward model, but we account for their effects on the radiances implicitly by deriving the MLER cloud fraction or the LER value and treating them as a function of wavelength using a second-order polynomial. The RRS spectrum ( $\sigma^{rrs}$ ), also known as the Ring cross section, is calculated using the method developed by Sioris and Evans [1999]. This method can take into account the sun-target-observer

geometry and the absorber (ozone and SO<sub>2</sub>) amounts, which are updated at each iteration step, accurately characterizing the spectral structure of RRS filling-in due to ozone and SO<sub>2</sub> absorption.

#### 4. Case Study: October 2005 Eruption of Sierra Negra

[10] A large eruption started at Sierra Negra volcano (Galápagos Islands) at about 5:30 pm local time on October 22, 2005 [Geist *et al.*, 2008], sending ash and steam clouds high into the atmosphere after the OMI overpass at 1:45 pm local time. These high clouds had drifted away by the next day when OMI made its first observation of this eruption, which became effusive with strong lava fountains (up to 300 m high) that produced long (>2 km) lava flows, manifested as the hot (orange) spot in Figure 1a, a composite Moderate Resolution Imaging Spectroradiometer (MODIS) image. This activity released very little ash but large amounts of gas, including SO<sub>2</sub>, into the lower troposphere, providing an excellent test of the ISF technique described in this paper.

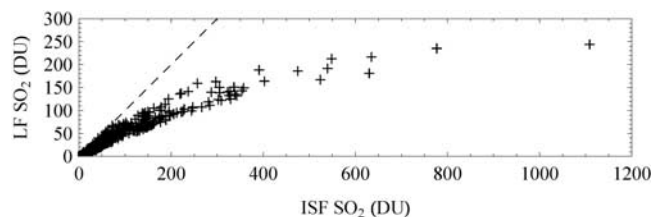
[11] We applied the ISF approach described above to the OMI measured spectra in the wavelength range 317.8 – 333 nm, which contains about 100 spectral measurements. This reduces the impact of random noise compared to using just a few discrete wavelengths selected from a full spectrum; consequently information contents can be more precisely retrieved. Since most of the measured TOA photons at these wavelengths have traversed the entire atmosphere, they provide information on total columns of the absorbers, but limited information about their vertical distributions. Therefore a priori profiles for both SO<sub>2</sub> and ozone are employed in our forward model. For ozone, we use the TOMS Version 8 (V8) profile climatology [Bhartia and Wellemeyer, 2002], which is uniquely determined by latitude, day of the year, and ozone amount. For the SO<sub>2</sub> plume, we represent its profile using a Gaussian, characterized by a central altitude and width.

[12] We demonstrate here that retrievals of very large total SO<sub>2</sub> columns are greatly improved using the new ISF algorithm. In Figure 2, we compare ISF ozone and SO<sub>2</sub> results with those from the operational OMI LF algorithm [Yang *et al.*, 2007] on October 23, 2005; both algorithms assumed a central plume altitude of ~2.5 km, which is appropriate for this case where SO<sub>2</sub> is mostly confined to the lower troposphere. Field observers on October 23 reported a plume ‘only several kilometres high’ [Geist *et al.*, 2008]. Figure 2a shows the ozone consistency achieved using the ISF method, with ISF ozone values unaffected by the presence of the SO<sub>2</sub> plume, whereas the operational LF ozone map (Figure 2b) contains artifacts, including high ozone at the plume edge and unrealistically low ozone in the interior of the plume. The internal consistency of the ISF ozone data serves as an indirect validation of the simultaneously retrieved SO<sub>2</sub> columns, which are generally larger than the corresponding LF columns. A slight reduction in retrieved ozone in the small area just over the volcanic vent implies that small ozone errors may be associated with retrievals of extremely large SO<sub>2</sub> columns. However, retrievals with ozone values preset to those outside the plume yield similar SO<sub>2</sub> columns, indicating that these ozone errors have no significant impact on the SO<sub>2</sub> results.

Figure 3 illustrates an approximate linear relationship between LF and ISF SO<sub>2</sub> columns below 100 DU, in agreement with the limitations of the LF method noted in our previous paper [Yang *et al.*, 2007]. Above 100 DU of SO<sub>2</sub>, the underestimation of SO<sub>2</sub> by the LF algorithm becomes more severe as the actual SO<sub>2</sub> column increases, and in these cases, ISF retrievals are needed to accurately retrieve the vertical columns. For the Sierra Negra eruption on October 23, 2005, the highest retrieved SO<sub>2</sub> column is located directly above the eruptive vent (see Figure 1b), with an ISF column exceeding 1100 DU. To our knowledge this is the largest SO<sub>2</sub> vertical column retrieved to date by a space borne instrument. For the same pixel, the LF algorithm yields 247 DU, a 78% underestimation. From OMI measurements on this day, the ISF SO<sub>2</sub> tonnage is ~1740 kilotons for the entire plume, nearly twice that (~950 kilotons) obtained using the LF data. The very high SO<sub>2</sub> column amounts directly over or close to the volcanic vent are probably the result of the high SO<sub>2</sub> emission rates associated with strong lava fountains and large lava flows, coupled with the small OMI footprint and slow dispersion of the volcanic plume due to light winds (wind speed ~3 m/s) at the time of the OMI observation.

## 5. Conclusions

[13] For strong volcanic eruptions that inject large amounts of SO<sub>2</sub> into the atmosphere, accurate quantification of the emissions from satellite BUV measurements requires proper accounting for the non-linear SO<sub>2</sub> absorption effect, and this can be accomplished using the ISF technique described in this paper. Application of the ISF method to



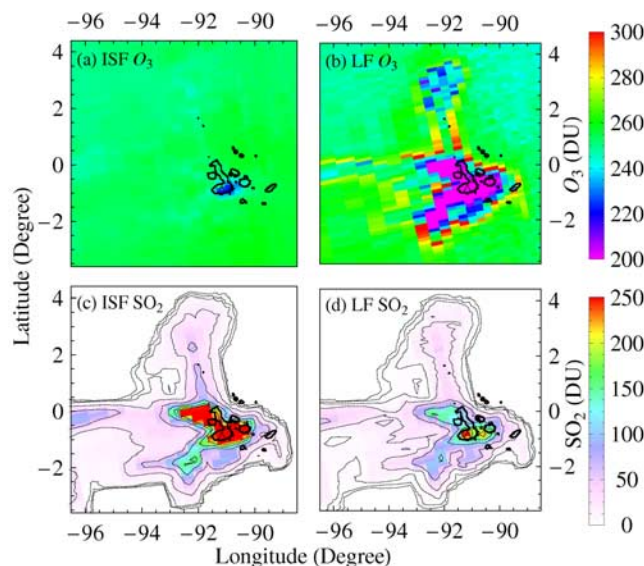
**Figure 3.** Comparison of ISF and LF SO<sub>2</sub> columns in the Sierra Negra eruption cloud on October 23, 2005. The LF retrievals saturate at about 200 DU in this case.

OMI observations of the Sierra Negra eruption on October 23, 2005 and comparison with the corresponding operational OMI SO<sub>2</sub> retrievals using the LF algorithm clearly demonstrate great improvements in quantification of large SO<sub>2</sub> columns above the LF upper valid range of ~100 DU. The iterative retrieval approach described here can be modified to use discrete wavelength measurements, such as those from the series of TOMS instruments, offering the potential to improve the long-term SO<sub>2</sub> emission record. Future work will include further improvement of the retrieval of volcanic SO<sub>2</sub> columns by deriving plume altitudes directly from BUV measurements.

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**Figure 2.** OMI ozone and SO<sub>2</sub> maps derived from OMI BUV radiances for the Sierra Negra volcanic cloud on October 23, 2005, using the ISF, (a) O<sub>3</sub> and (c) SO<sub>2</sub>, and LF, (b) O<sub>3</sub> and (d) SO<sub>2</sub>, algorithms assuming a fixed SO<sub>2</sub> plume central altitude of 2.5 km and a plume width of 3.5 km (defined as the full width at half maximum of the Gaussian profile) for ISF and a width of ~5 km (i.e., the width of Umkehr layer 0) for LF.

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